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**Synthesis, crystal structure, and electrical and magnetic properties of
BaMo₆Te₆: a novel reduced molybdenum telluride containing infinite
chains of trans-face shared Mo₆ octahedra**

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Abstract

Powder samples and single crystals of the new ternary compound BaMo_6Te_6 were obtained by solid state reaction. The structure was determined by single-crystal X-ray diffraction. BaMo_6Te_6 crystallizes in the hexagonal space group $P6_3/m$ (No. 176) with unit-cell parameters $a = 9.3941(2) \text{ \AA}$, $c = 4.5848(1) \text{ \AA}$ and $Z = 1$. Full-matrix least-squares refinement on F^2 using 452 independent reflections for 17 refinable parameters resulted in $R_1 = 0.0208$ and $wR_2 = 0.0539$. The structure consists of one-dimensional infinite chains of trans-face shared Mo_6 octahedra capped by Se atoms. These chains that are running along the c axis are separated from each other by nine-coordinate Ba atoms. Resistivity measurements on a single crystal indicated that the BaMo_6Te_6 compound is metallic down to 160 K and semiconductor below. Magnetic susceptibility measurements showed that BaMo_6Te_6 is weakly diamagnetic with no anomaly at the metal-semiconductor transition.

Keywords

Reduced molybdenum chalcogenide, infinite molybdenum chains, magnetic susceptibility, resistivity measurement.

INTRODUCTION

Since the discovery of the Chevrel phases MMo_6X_8 (M = cation; X = S, Se, Te) [1] containing octahedral Mo_6 clusters about forty years ago, a plethora of ternary and quaternary reduced molybdenum chalcogenides containing larger clusters have been synthesized. The latter compounds contain clusters such as Mo_9 , Mo_{12} , Mo_{15} , Mo_{18} , Mo_{21} , Mo_{24} , Mo_{30} and Mo_{36} . All of these clusters result from the uniaxial trans-face-condensation of octahedral Mo_6 clusters and are present, for example, in the series of compounds $\text{M}_{n-2}\text{Mo}_{3n}\text{X}_{3n+2}$ (M = Rb, Cs; X = S, Se or Te; n = 3, 4, 5, 6, 7, 8, 10 and 12) [2-12]. The members with even n crystallize in the rhombohedral trigonal space group $R\bar{3}$ with a structure closely related to the ternary chalcogenides MMo_6X_8 and the members with odd n in the hexagonal space group $P6_3/m$. The final step of the condensation corresponds to the infinite chain of trans-face-sharing Mo_6 octahedra $\left| \text{Mo}_{6/2} \right|_{\infty}^1$ [13, 14] found in the hexagonal compounds $\text{M}_2\text{Mo}_6\text{X}_6$ (M = Na, K, Rb, Cs, In, Tl; X = S, Se, Te) and monoclinic one, AgMo_6Te_6 [15]. The interest for these compounds lies not only in their fascinating structural aspect but also in their interesting physical properties. Indeed, the sulfides and selenides generally present superconducting or metal-insulator transitions at low temperature. Thus, studies of the normal and superconducting states of both $\text{Cs}_2\text{Mo}_{12}\text{Se}_{14}$ (n = 4) and $\text{Rb}_4\text{Mo}_{18}\text{Se}_{20}$ (n = 6) by measuring the conductivity and magnetization of single-crystals and powder samples have shown that these compounds can be classified among the "exotic" superconductors [16]. The quasi 1-D superconductor $\text{Tl}_2\text{Mo}_6\text{Se}_6$ presents extreme type II and non-B.C.S. behaviors [17-20]. In addition, the anisotropy of the electronic properties in the latter compound is one of the largest ever observed in a superconductor with the ratio of the conductivities parallel and perpendicular to the infinite chains $(\sigma_{//}/\sigma_{\perp})$ of about 1000, and the ratio of the upper critical fields $(H_{C2//}/H_{C2\perp})$ of about 26. In contrast, the $\text{M}_2\text{Mo}_6\text{X}_6$ compounds with M being an element of the group IA present a continuous metal-nonmetal transition between 100 and 150 K [21]. Recent local density approximation calculations show that in all members of the $\text{M}_2\text{Mo}_6\text{Se}_6$ series, a single 1D helix band crosses the Fermi level. Its in-plane dispersion is reduced by a factor of ten in $\text{Rb}_2\text{Mo}_6\text{Se}_6$ compared with $\text{Tl}_2\text{Mo}_6\text{Se}_6$ and $\text{In}_2\text{Mo}_6\text{Se}_6$ while the out-of-plane dispersion is practically unchanged [22]. This reduction is sufficient to explain the trend from metallic conductivity followed by strongly-coupled superconductivity in $\text{Tl}_2\text{Mo}_6\text{Se}_6$ to a high-temperature metal-insulation transition in $\text{Rb}_2\text{Mo}_6\text{Se}_6$. This insulating

behaviour is consequence of either a Peierls distortion (Charge Density Wave or CDW), a Spin Density Wave or an Anderson localization. In 1985, Chevrel et al. [23] mentioned the existence of an isomorphous compound, $\text{Ba}_2\text{Mo}_6\text{Te}_6$. However, such a stoichiometry would imply 14 electrons per Mo_3Te_3 fragment and thus, which seems improbable, the occupation of antibonding states since the optimal count per Mo_3Te_3 is 13 [24]. Thus, we decided to reinvestigate this compound by growing single crystals, in particular. This study has led that the right stoichiometry is in fact BaMo_6Te_6 . We describe the synthesis, crystal structure, and electrical and magnetic properties of this compound here.

EXPERIMENTAL,

Synthesis and Crystal Growth

X-ray diffractometrically pure powder of BaMo_6Te_6 was prepared from a stoichiometric mixture of BaTe, MoTe_2 and Mo. BaTe and MoTe_2 were prepared, respectively, by heating the elements together in silica tubes at 800 °C during 48h. Before use, the Mo powder was heated under a hydrogen flow at 1000°C for 6 hours. The starting reactants were then pressed into pellets (ca. 5g) and loaded into a molybdenum crucible which was previously outgassed at 1500 °C for 15 min under a dynamic vacuum of about 10^{-5} Torr. The Mo crucible was subsequently sealed under a low argon pressure using an arc-welding system. The sample was heated at a rate of 300°C/hour to 1400 C, kept at the temperature for 48 hours and then cooled at 100°C/hour down to 1100°C at which point the furnace was shut down and allowed to cool to room temperature. The resulting product was found to be single-phase on the basis of its X-ray powder diffraction pattern carried out on a D8 Bruker Advance diffractometer equipped with a LynxEye detector ($\text{CuK}\alpha_1$ radiation) and operating at 40 kV and 40 mA. (Figure 1). Crystals were grown by heating a mixture of composition “ $\text{Ba}_2\text{Mo}_6\text{Te}_6$ ” in a sealed molybdenum crucible at about 1500 °C for 48h. The crucible was then cooled at a rate of 100 °C/hour down to 1000 °C, and finally furnace-cooled to room temperature. Crystals thus obtained have the shape of needles with hexagonal cross section.

Single Crystal X-ray Study

The X-ray diffraction data were collected on a Nonius Kappa CCD diffractometer using graphite-monochromated Mo-K α radiation ($\lambda = 0.71073 \text{ \AA}$) at room temperature. The COLLECT program package [25] was employed to establish the angular scan conditions (φ and ω scans) used in the data collection. The data set was processed using EvalCCD [26] for the integration procedure. An absorption correction was applied using the description of the crystal faces and the analytical method described by de Meulenaar and Tompa [27]. Analysis of the data revealed that the systematic absence $(000l) \quad l = 2n+1$ was only consistent with the hexagonal space group $P6_3/m$. The structure was solved with the direct methods program SIR97 [28] and refined using SHELXL97 [29] in the $P6_3/m$ space group. Refinement of the occupancy factor of the Ba site led to a value of 0.492(3) indicated a half occupation of the cationic site compared to the other hexagonal compounds $M_2Mo_6X_6$ ($M = Na, K, Rb, Cs, In, Tl$; $X = S, Se, Te$). An attempt to refine the structure in the acentric space group $P6_3$ was unsuccessful and led to a higher R factor as well as larger residual electron density peaks. Because of the disordering of the Ba atom, we made long-exposure rotation photographs along the a and c axes on a single crystal of $BaMo_6Te_6$. The latter did not reveal any superlattice reflection. Crystallographic data and X-ray structural analysis for the $BaMo_6Te_6$ compound are summarized in Table 1. The final atomic coordinates, and the equivalent isotropic displacement parameters are gathered in Table 2, and selected interatomic distances are listed in Table 3. Further details of the crystal structure investigation can be obtained from the Fachinformationszentrum Karlsruhe, 76344 Eggenstein-Leopoldshafen, Germany, (fax: (49) 7247-808-666; e-mail: crysdata@fiz.karlsruhe.de) on quoting the depository number CSD- 427614.

Electrical Resistivity Measurements

The study of the temperature dependence of the electrical resistivity was carried out on a single crystal of $BaMo_6Te_6$ using a conventional ac four-probe method with a current amplitude of 0.1 mA. Contacts were ultrasonically made with molten indium on the single crystal previously characterized on a Nonius Kappa CCD diffractometer. The ohmic behavior and the invariance of the phase were checked during the different measurements at low and room temperature.

Magnetic susceptibility measurements.

Susceptibility data were collected on cold pressed powder samples (ca. 100 mg) using a Quantum Design SQUID magnetometer between 4.2 K and 400 K and at an applied field of 0.1 T.

RESULTS AND DISCUSSION

Crystal Structure

The comparison of the unit cell parameters of the previous compound reported as “Ba₂Mo₆Te₆” ($a = 9.41 \text{ \AA}$, $c = 4.57 \text{ \AA}$, $V = 350 \text{ \AA}^3$) with those of the title compound ($a = 9.3941(2) \text{ \AA}$, $c = 4.5848(1) \text{ \AA}$, $V = 350.40(1) \text{ \AA}^3$) clearly shows that we are in the presence of the same phase i.e. BaMo₆Te₆.

A view of the unit cell of BaMo₆Te₆ is given in Figure 1. The structure is composed of infinite chains of trans-face-sharing Mo₆ octahedra $\left| \text{Mo}_{6/2} \right|_{\infty}^1$ capped by Te atoms. These Mo-Te chains run along the c axis and have crystallographically imposed symmetry -3 . The Ba²⁺ cations separate the $\left| \text{Mo}_{6/2} \text{Te}_{6/2} \right|_{\infty}^1$ chains and are each coordinated to nine Te atoms in a tricapped trigonal-prismatic arrangement with imposed -6 symmetry. The Mo-Mo distances within and between the Mo₃ triangles perpendicular to the c axis are 2.6602(6) and 2.7593(2) Å, respectively. The latter values are of the same order than those observed in other reduced molybdenum tellurides such as Mo₆Te₈, Cs₂Mo₁₂Te₁₄, Cs₃Mo₁₅Te₁₇, Cs₄Mo₁₈Te₂₀ and Cs₆Mo₂₄Te₂₆ [5] in which the average intratriangle distances range from 2.663 to 2.688 Å and the average intertriangle ones from 2.765 to 2.785 Å. However, they differ significantly from those found in In₂Mo₆Te₆ (2.639 and 2.755 Å) and Tl₂Mo₆Te₆ (2.619 and 2.746 Å), the structures of which have been solved from RX powder data [15]. The Mo-Te distances range from 2.8015(4) to 2.8560(2) Å as usually observed. The Ba-centered tricapped trigonal prisms share triangular faces along [001]. The distance between the Ba and the Te atoms forming the prism is 3.4799(2) Å while that with the Te atoms capping the prism is longer, 3.6139(3) Å. These distances are in agreement with that found in BaTe (3.502 Å), for example.

Electrical and Magnetic Properties.

Electrical resistivity measurements were performed on a single-crystal of BaMo₆Te₆. The temperature dependence of the electrical resistivity of BaMo₆Te₆ measured along the

chain growth is similar to that reported for the $M_2Mo_6X_6$ compounds ($M = Na, K, Rb$ and Cs ; $X = Se, Te$) (Fig. 5). Thus, $BaMo_6Te_6$ is metallic in the temperature range 160–290 K with a resistivity of about 0.038 m Ω .cm at room temperature and 0.022 at 160 K. Below 160 K, a semiconductor behaviour is observed since the electrical resistivity increases with decreasing temperature. The activation energy deduced from the linear part of the $1/T$ vs. $\log\rho$ plot is 220 K that is close to that reported for the $Cs_2Mo_6Te_6$ compound (ie 225 K) [21].

The temperature dependence of the magnetic susceptibility of $BaMo_6Te_6$ is shown in Figure 5. The susceptibility of the $BaMo_6Te_6$ compound is nearly temperature-independent and weakly diamagnetic in the range 100–300 K with a value of about -7×10^{-8} emu/g closed to that reported for some $M_2Mo_6X_6$ compounds with M being an element of the group IA [21]. This behavior is consistent with the absence of localized moments on the Mo atom network. The low-temperature upturn could be attributed to small amounts of paramagnetic impurities often present in the starting reactants. It is interesting to note that no anomaly was observed at the metal-nonmetal transition as previously observed for the $M_2Mo_6X_6$ compounds [21]. This is reminiscent of the well known Anderson localisation that arises from defects [30]. Indeed, unlike the CDW transitions [31, 32] that require an opening of gap near the Fermi level and associated changes in the structural and variable temperature magnetization studies, the present compound does not show any reduction in the magnetization value near the metal-nonmetal transition. In addition, the broad nature of the transition in the electrical resistivity also supports defect-induced localization of the carriers. A similar observation has also been reported for KMo_4O_6 that contains infinite chains of trans-edge shared Mo_6 octahedral clusters [33].

CONCLUSION

In summary, the new compound $BaMo_6Te_6$ has been synthesized by solid state reaction. This compound is isomorphous with the $M_2Mo_6X_6$ ($M = Na, K, Rb, Cs, In, Tl$; $X = S, Se, Te$) compounds and, thus its crystal structure contains infinite chains of trans-face-sharing Mo_6 octahedra $\left| Mo_{6/2} \right|_{\infty}^1$ capped by Te atoms. The Ba^{2+} cations occupy statically half of the tricapped trigonal prisms formed by the Te atoms between the chains while, in the $M_2Mo_6X_6$ compounds, all the tricapped trigonal prismatic sites are occupied. The temperature dependences of the electrical resistivity and magnetic susceptibility are similar to those observed for the $M_2Mo_6X_6$ compounds with M being an element of the group IA. In

particular, the broad metal-nonmetal transition around 160 K results probably from an Anderson type localization due to the absence of anomaly in the magnetization curve. Finally, it would be interesting to study the thermoelectric properties of BaMo_6Te_6 and $\text{M}_2\text{Mo}_6\text{X}_6$ ($\text{M} = \text{Na}, \text{K}, \text{Rb}, \text{Cs}$; $\text{X} = \text{S}, \text{Se}, \text{Te}$). Indeed, recent works on the finite chain cluster compounds $\text{Ag}_x\text{Mo}_9\text{Se}_{11}$ ($3.4 \leq x \leq 3.9$) and $\text{Ag}_3\text{In}_2\text{Mo}_{15}\text{Se}_{19}$ containing Mo_9 clusters [34-36] have shown that Mo cluster chalcogenides are good candidates for thermoelectricity.

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Figure Captions

Figure 1. Observed (crossed), calculated (black line) and difference profiles for the refinement of BaMo₆Te₆ in profile-matching mode ($\lambda = 1.5406 \text{ \AA}$).

Figure 2. The crystal structure of BaMo₆Te₆ as viewed down the c axis, parallel to the direction of the Mo_{6/2} chain growth. Ellipsoids are drawn at the 97 % probability level.

Figure 3. The crystal structure of BaMo₆Te₆ as viewed perpendicularly to the Mo_{6/2} chain growth (ellipsoids at the 97 % probability level).

Figure 4. Tellurium environment for the Ba atom (ellipsoids at the 97 % probability level).

Figure 5. Temperature dependence of the resistivity for BaMo₆Te₆. The inset shows the Arrhenius plots.

Figure 6. Magnetic susceptibility of BaMo₆Te₆ as a function of temperature. Data were taken under an applied field of 0.1 T.

Table 1. Crystal data and structure refinements of BaMo₆Te₆

Empirical formula	BaMo ₆ Te ₆
Formula weight (g mol ⁻¹)	1478.58
Crystal system, space group	hexagonal, P6 ₃ /m
Unit cell dimensions (Å, deg)	$a = 9.3941$ (2), $c = 4.5848$ (1)
Volume (Å ³)	350.40 (1)
Z, Calculated density (g/cm ³)	1, 7.007
Absorption coefficient (mm ⁻¹)	20.151
Crystal color and habit	Black, needle like
Crystal size (mm ³)	0.152 × 0.063 × 0.057
Theta range for data collection (deg)	2.50–32
Limiting indices	$-14 \leq h \leq 14$, $-14 \leq k \leq 14$, $-5 \leq l \leq 6$
Reflections collected/unique	8209/452
R (int)	0.0444
Absorption correction	Analytical
Max./min. transmission	0.4142/0.1328
Data/restraints/parameters	3903/0/101
Goodness-of-fit on F^2	1.081
R indices [$I > 2\sigma(I)$]	$R_1=0.0208$, $wR_2=0.0539$
Extinction coefficient	0.00418(11)
Largest diff. peak and hole (eÅ ⁻³)	1.437 and -1.177

Table 2. Atomic coordinates and equivalent isotropic displacement parameters (\AA^2) for the compound BaMo_6Te_6 .

Atom	Wyckoff position	x	y	z	Ueq	s.o.f
Mo	6h	-0.14384(4)	0.03396(4)	0.7500	0.00504(14)	1
Te	6h	0.07032(3)	0.37457(3)	0.7500	0.00659(13)	1
Ba	2c	0.3333	0.6667	0.2500	0.0075(3)	0.492(3)

Table 3. Selected Interatomic Distances (Å) for BaMo₆Te₆.

Mo-Mo (X2)	2.6602(6)	Ba-Te (X6)	3.4799(2)
Mo-Mo (X4)	2.7593(2)	Ba-Te (X3)	3.6139(3)
Mo-Te	2.8015(4)		
Mo-Te	2.8117(4)		
Mo-Te (X2)	2.8560(2)		

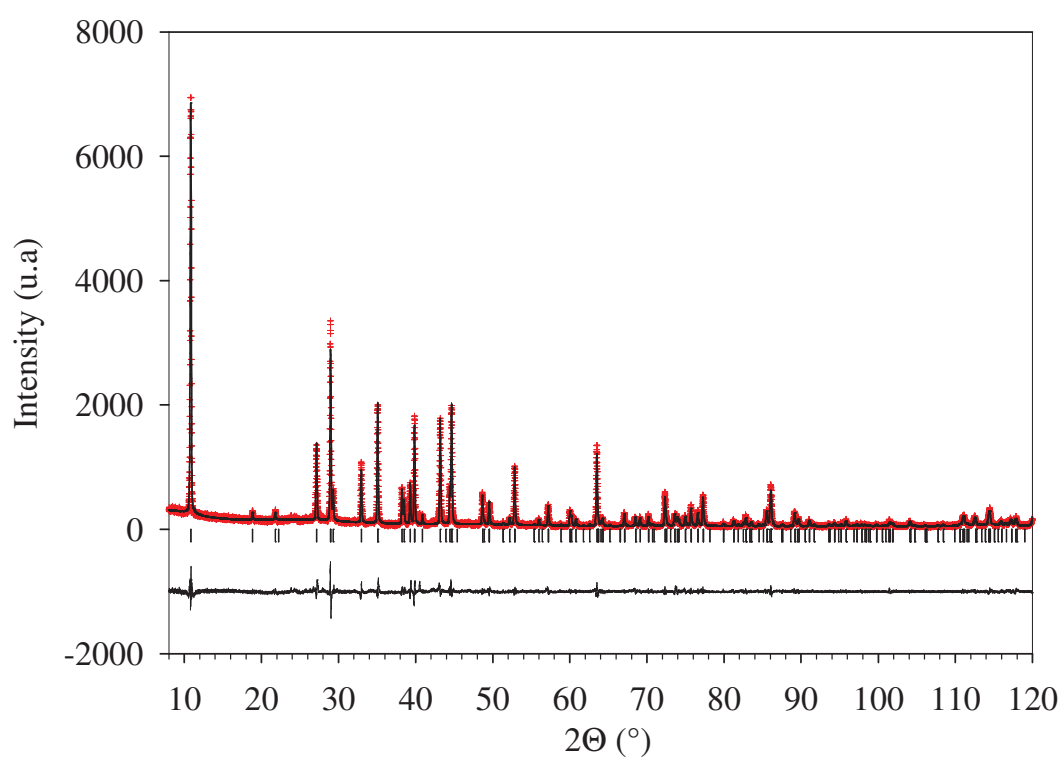


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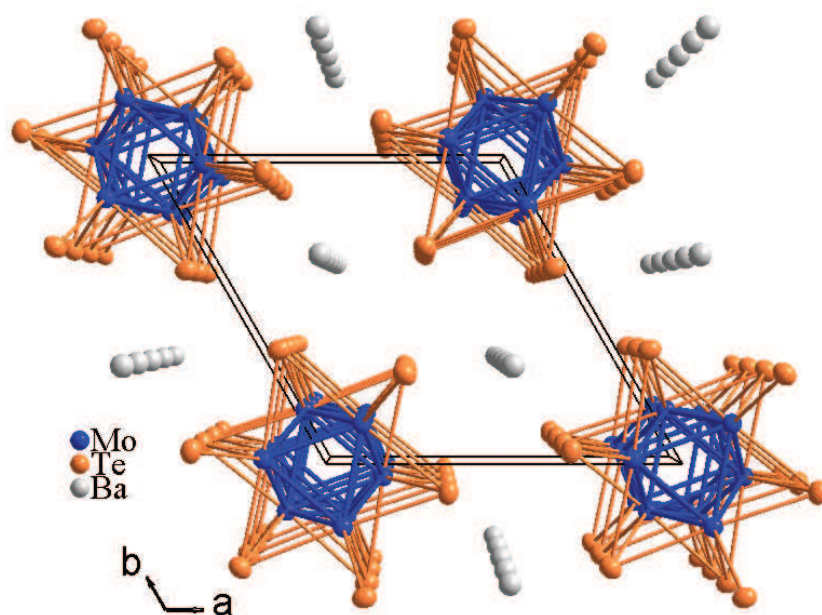


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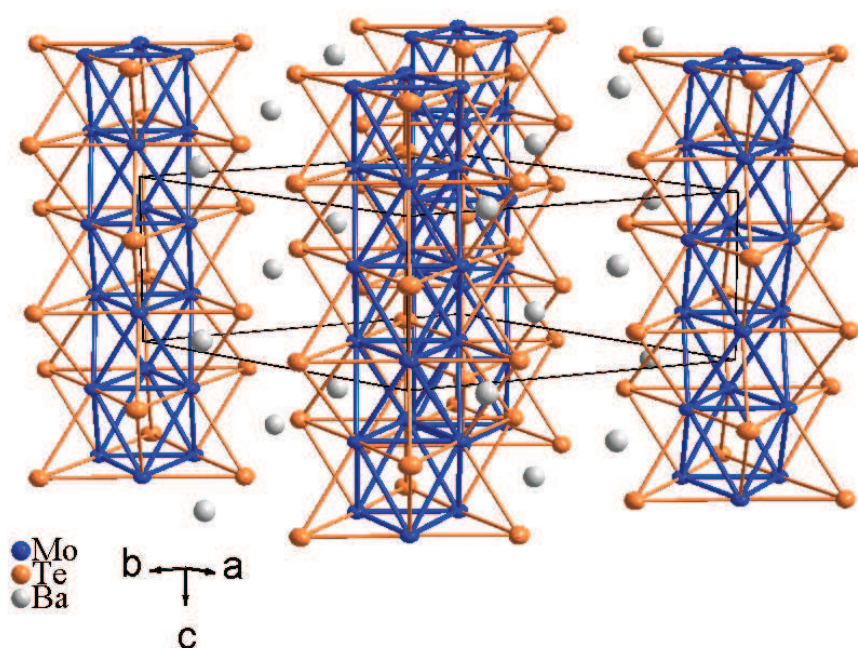


Figure 3. The crystal structure of BaMo_6Te_6 as viewed perpendicularly to the $\text{Mo}_{6/2}$ chain growth (ellipsoids at the 97 % probability level).

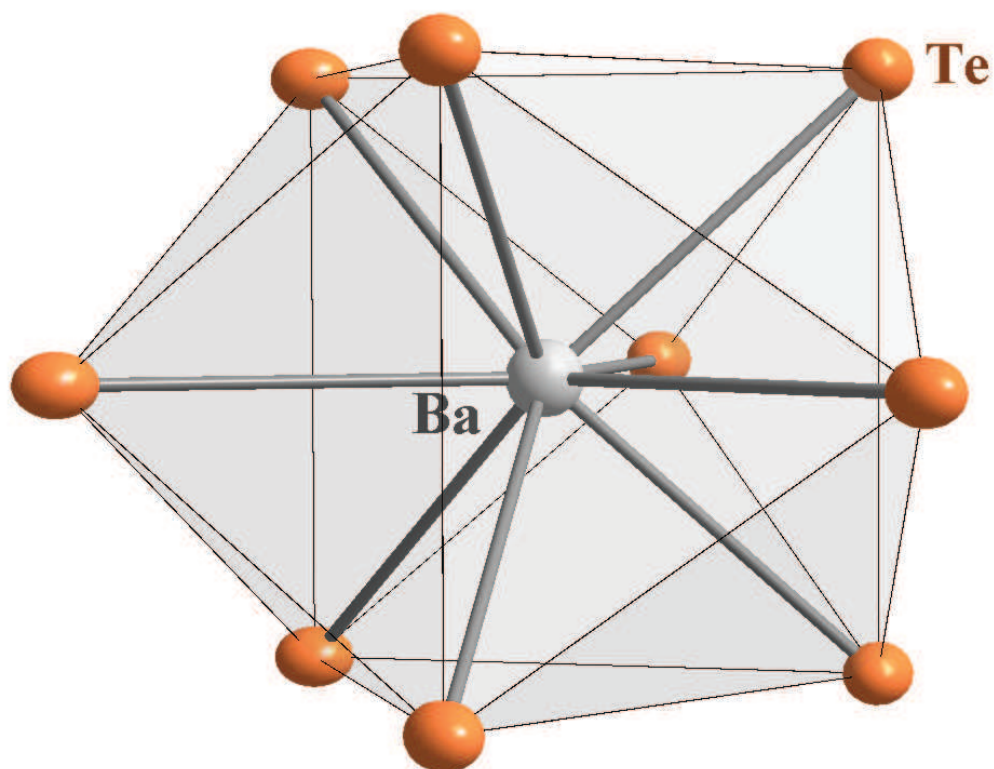


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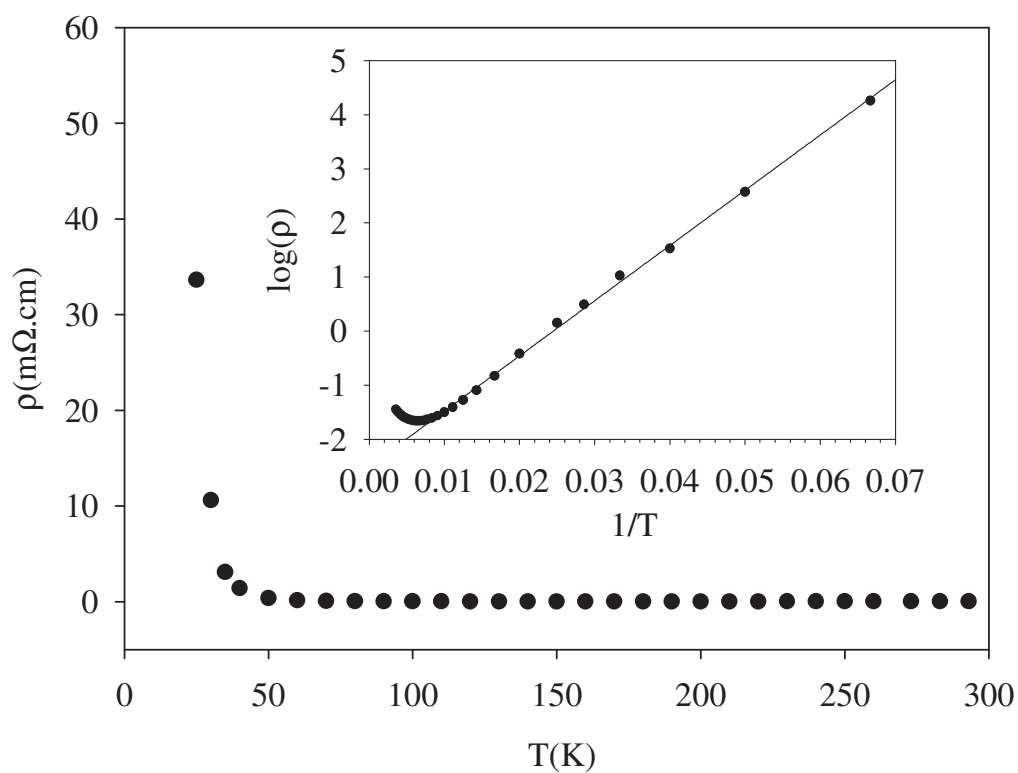


Figure 5. Temperature dependence of the resistivity for BaMo_6Te_6 . The inset shows the Arrhenius plots.

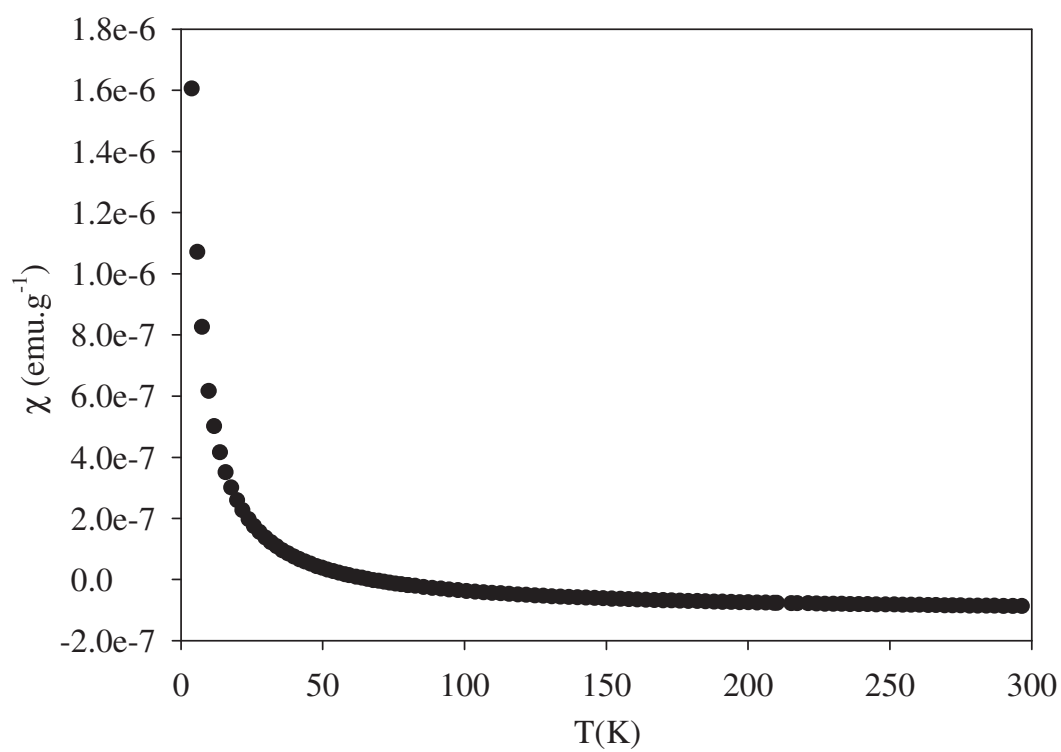


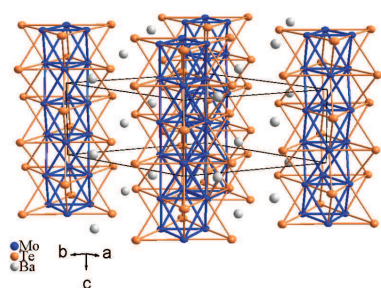
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We present here the synthesis, the crystal structure, and the electrical and magnetic properties of the new compound BaMo_6Te_6 containing infinite chains of trans-face shared Mo_6 octahedra



Highlights

- . BaMo₆Te₆ contains infinite chains of trans-face-sharing Mo₆ octahedra $\left| \text{Mo}_{6/2} \right|_{\infty}^1$.
- . Synthesis by solid state reaction.
- . Single-crystal X-ray study.
- . Continuous metal-nonmetal transition.
- . Anderson localization.